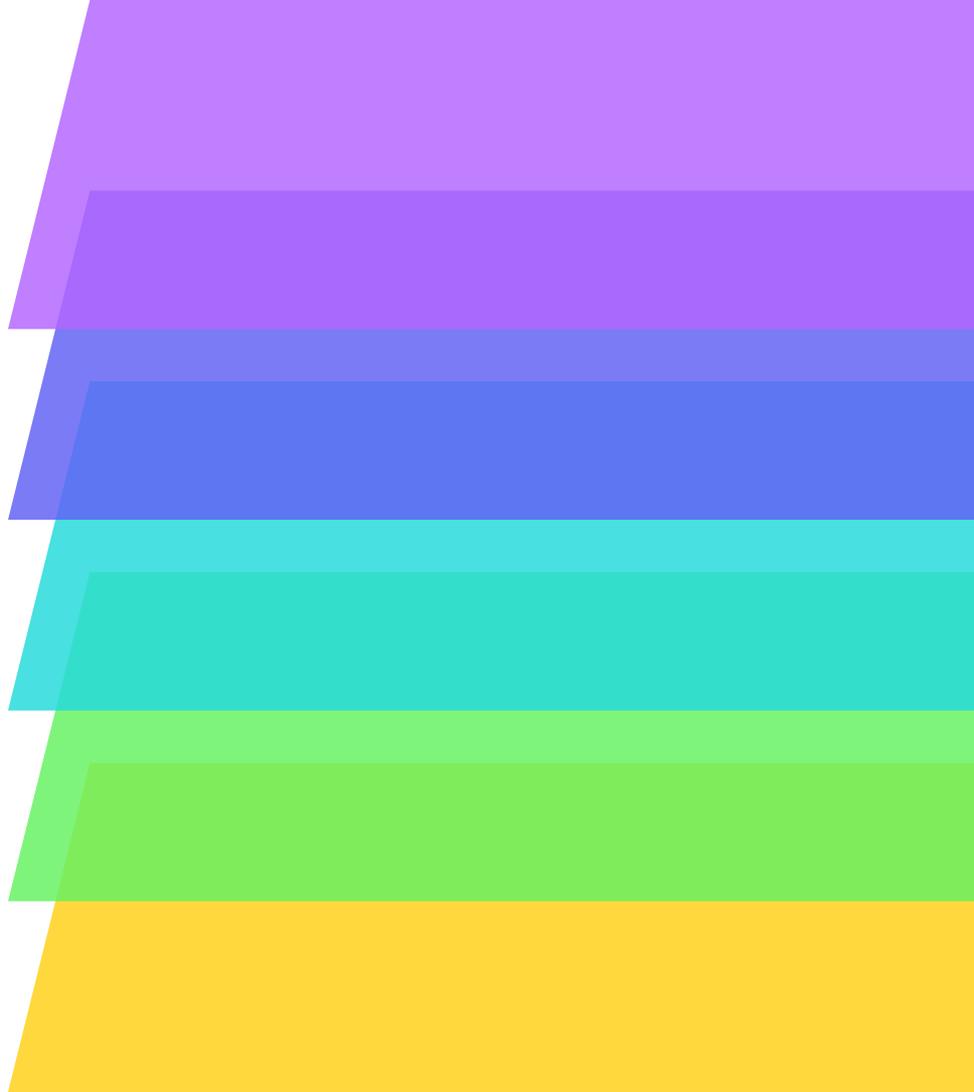


Magnetic monopoles

Diego França de Oliveira



1ST

Predicting Magnetic Monopoles

Quantization of charge, gauge symmetries and super massive monopoles

2ND

Looking for magnetic monopoles

Detection experiments, astronomical bounds and the “monopole problem”

3RD

Creating a magnetic monopole

MoEDAL experiment

1ST

Predicting Magnetic Monopoles

Maxwell's equation symmetry (in gaussian units)

In vacuum

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 0 \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= 0\end{aligned}$$

With charge

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 4\pi\rho \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= \frac{4\pi}{c} \vec{J}.\end{aligned}$$

Maxwell's equation symmetry (in gaussian units)

In vacuum

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial B}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 0 \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial E}{\partial t} &= 0\end{aligned}$$

With charge

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 4\pi\rho_m \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial B}{\partial t} &= -\frac{4\pi}{c} \vec{J}_m \\ \vec{\nabla} \cdot \vec{E} &= 4\pi\rho \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial E}{\partial t} &= \frac{4\pi}{c} \vec{J}.\end{aligned}$$

The absence of magnetic monopoles

Electrodynamics

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$\nabla \cdot \vec{B} = \nabla \cdot (\vec{\nabla} \times \vec{A})$$

$$\nabla \cdot \vec{B} = 0$$

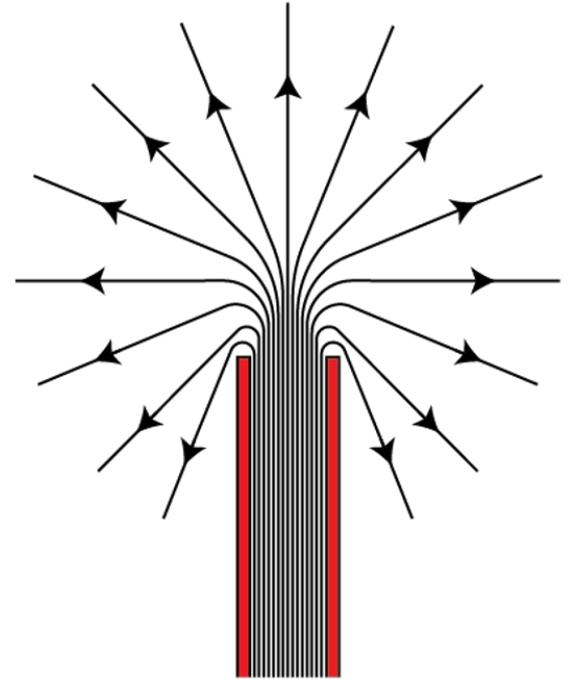
Quantum mechanics

$$|\psi_C \rangle = \exp\left(\frac{-iq}{\hbar c} \int \vec{A} \cdot d\vec{l}\right) |\psi_A \rangle$$

Potential fields forbid magnetic monopoles in ED and QM

Dirac's argument (1931)

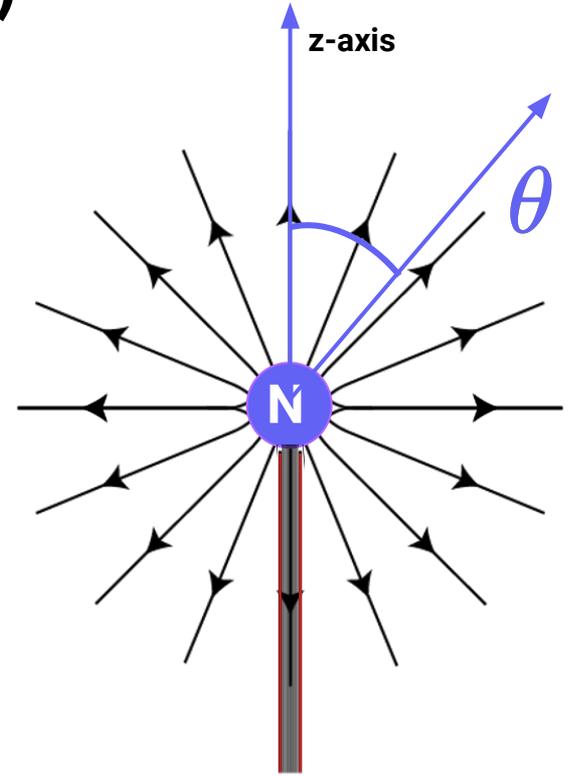
- A current through a solenoid produces magnetic field
- If the solenoid is infinitely long and thin, the field lines never close
- We have something similar to a magnetic monopole, but does the theory hold up?



Dirac's argument (1931)

Magnetic field produced: $\vec{B} = \frac{g}{r^2} \hat{r}$

A possible vector potential: $\vec{A} = g \frac{1 - \cos(\theta)}{r \sin(\theta)} \hat{\varphi}$

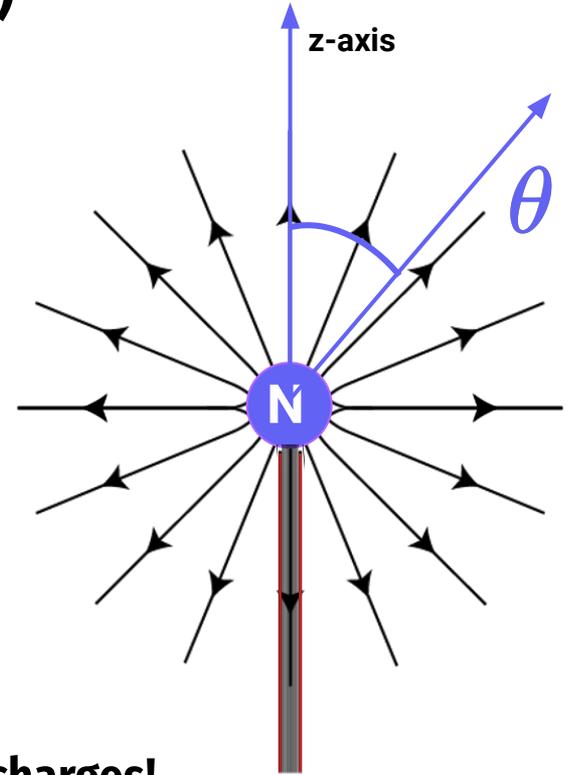


Dirac's argument (1931)

Flux of magnetic field:

$$\begin{aligned}\Phi_B &= \int_A \vec{B} \cdot d\vec{S} = \int_\Gamma \vec{A} \cdot d\vec{l} \\ &= 2\pi g(1 - \cos(\theta))\end{aligned}$$

$$\Phi_B(\theta = 2\pi) = 4\pi g$$



Similar to the gauss law for electric charges!

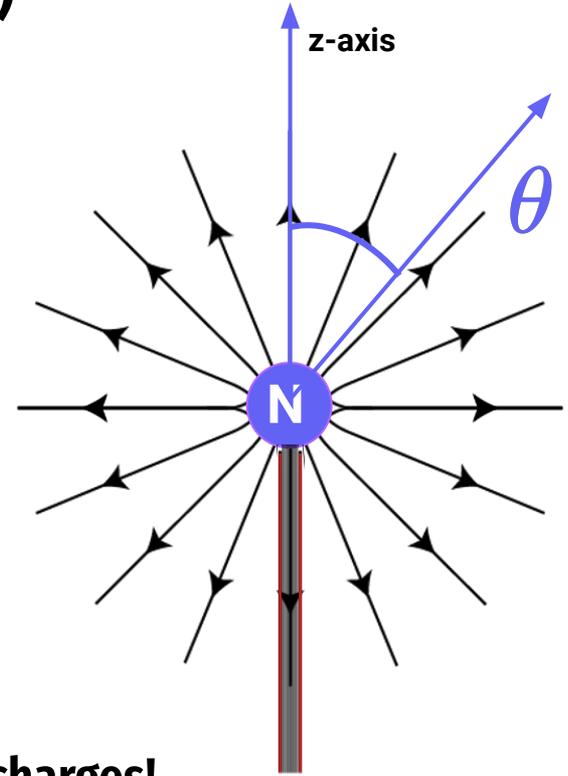
$$\Phi_E(\theta = 2\pi) = 4\pi q$$

Dirac's argument (1931)

Flux of magnetic field:

$$\begin{aligned}\Phi_B &= \int_A \vec{B} \cdot d\vec{S} = \int_\Gamma \vec{A} \cdot d\vec{l} \\ &= 2\pi g(1 - \cos(\theta))\end{aligned}$$

$$\Phi_B(\theta = 2\pi) = 4\pi g$$



Similar to the gauss law for electric charges!
But is mathematically wrong

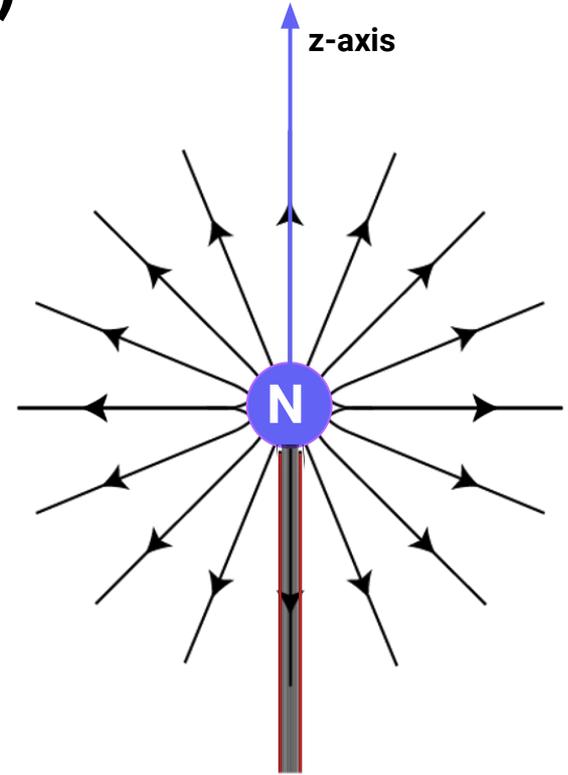
Dirac's argument (1931)

Flux of magnetic field:

$$\Phi_B = \begin{cases} 2\pi g(1 - \cos(\theta)); & \theta < \pi \\ 0 & ; \theta = \pi \end{cases}$$

$$\vec{B} = \nabla \times \vec{A} - 4\pi g \Theta(-z) \delta(x) \delta(y) \hat{z}$$

Dirac string singularity
Physical or mathematical?



Dirac's argument (1931)

Is the dirac string detectable?

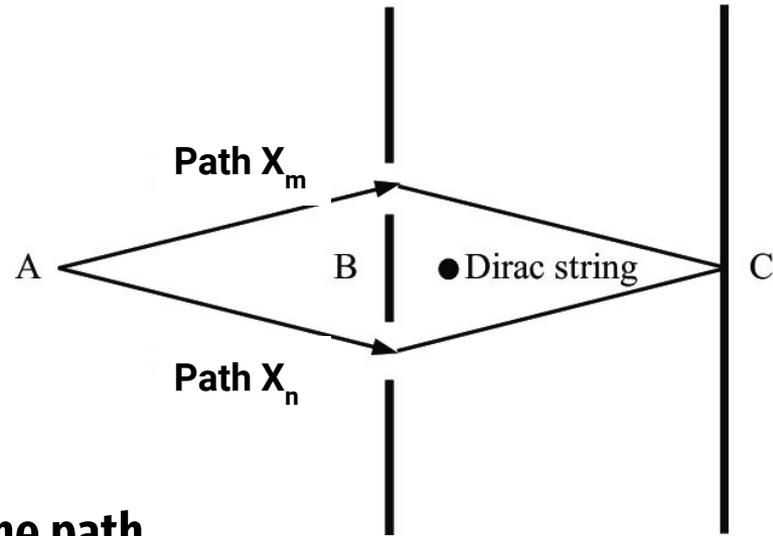
By the Aharonov-bohm effect:

$$|\psi_C \rangle = \exp\left(\frac{-iq}{\hbar c} \int \vec{A} \cdot d\vec{l}\right) |\psi_A \rangle$$

The phase change will be

$$\theta = \frac{q}{\hbar c} \int \vec{A} \cdot d\vec{l}$$

Measuring the phase in C allows one to know the path.



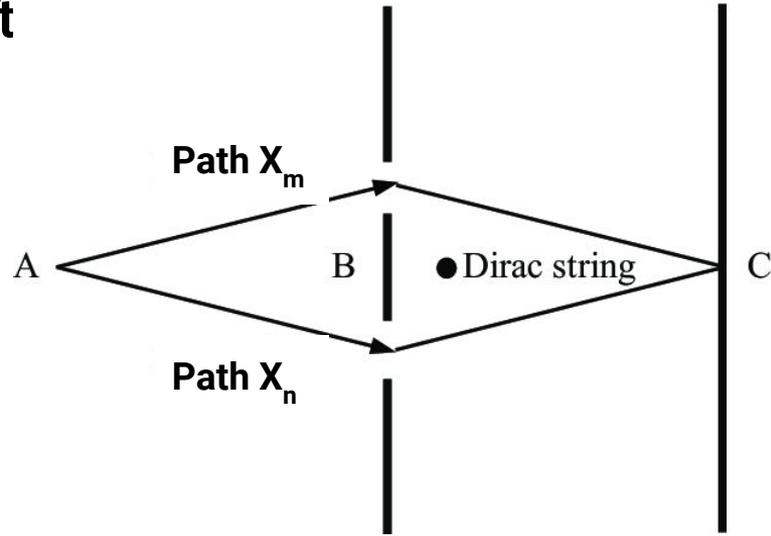
Dirac's argument (1931)

It's possible to calculate the difference in phase shift
Between paths.

$$\Delta\theta = \frac{-q}{\hbar c} \int_{A \rightarrow X_m \rightarrow C} \vec{A} \cdot d\vec{l} - \frac{-q}{\hbar c} \int_{A \rightarrow X_n \rightarrow C} \vec{A} \cdot d\vec{l}$$

$$\Delta\theta = \frac{-q}{\hbar c} \int_A \vec{B} \cdot d\vec{S}$$

$$\Delta\theta = \frac{-q}{\hbar c} 4\pi g$$



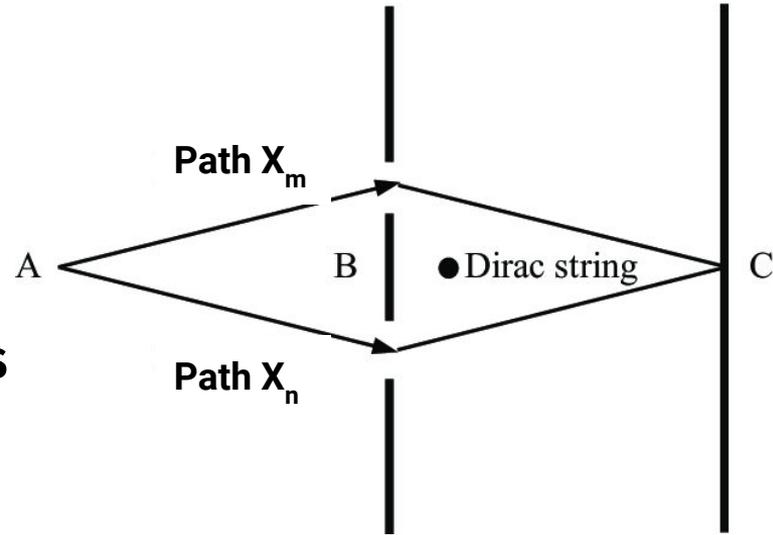
Dirac's argument (1931)

There is a condition where is undetectable:

$$\frac{-q}{\hbar c} 4\pi g = n2\pi$$

$$\frac{-qg}{\hbar c} \in \mathbf{Z}$$

If **charge is quantized**, then magnetic monopoles are allowed!

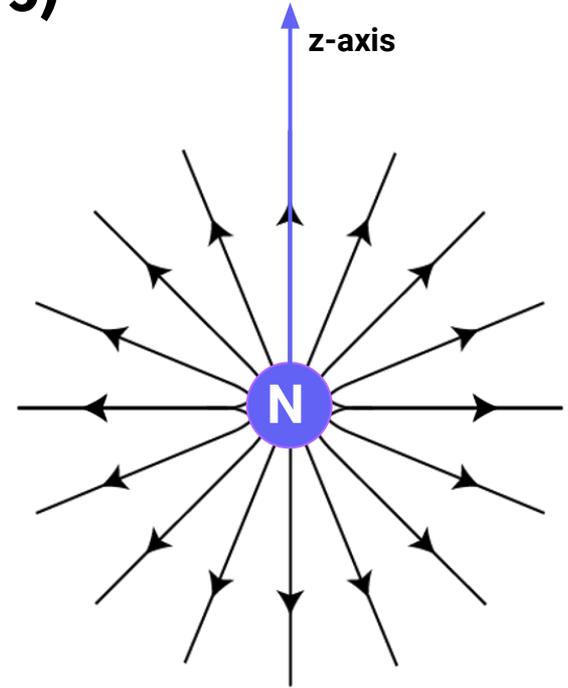


Removing the string (1975)

Charge quantization comes from the dirac string situation. Can we remove the string while getting the quantization condition?

$$\nabla \cdot \vec{B} = \nabla \cdot (\vec{\nabla} \times \vec{A})$$

Not using a vector potential...Let's try to use more



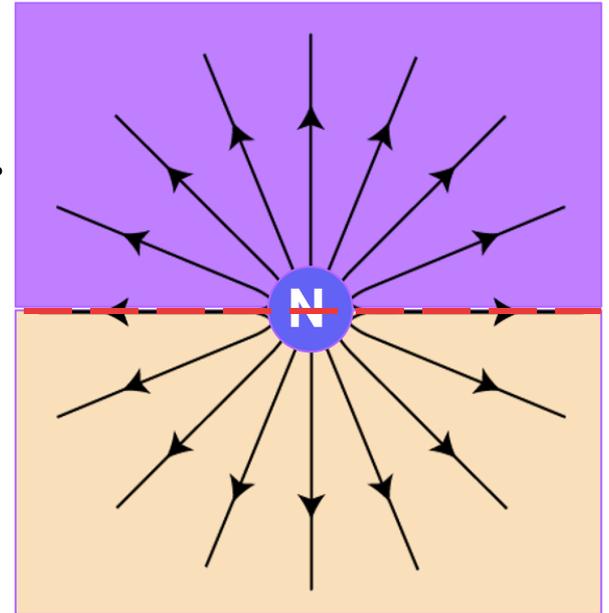
Removing the string (1975)

Let's assume two vector potential in two different regions.

$$\vec{A}_{sup} = g \frac{1 - \cos(\theta)}{r \sin(\theta)} \hat{\varphi}; 0 \leq \theta \leq \frac{\pi}{2}$$

$$\vec{A}_{inf} = -g \frac{1 + \cos(\theta)}{r \sin(\theta)} \hat{\varphi}; \frac{\pi}{2} \leq \theta \leq \pi$$

$$\left. \begin{array}{l} \vec{A}_{sup} \\ \vec{A}_{inf} \end{array} \right\} \vec{B} = \frac{g}{r^2} \hat{r}$$



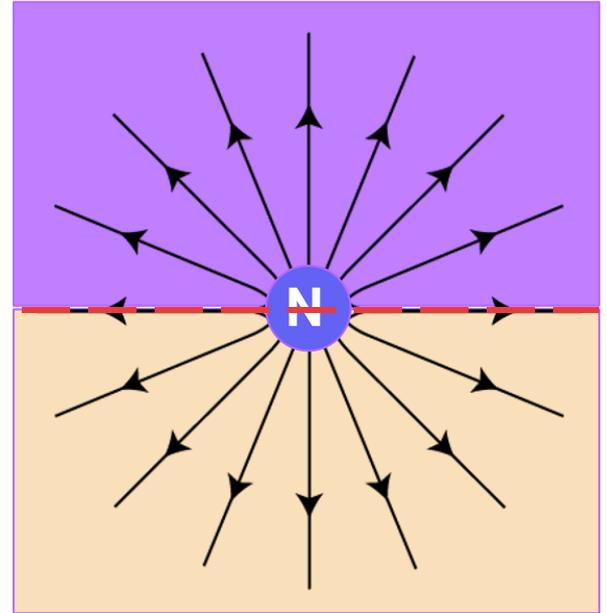
Removing the string (1975)

In the equator, is mandatory that:

$$\vec{\nabla}\chi = \vec{A}_{sup} - \vec{A}_{inf}$$

$$\vec{\nabla}\chi = -\frac{2g}{r\sin(\theta)}\hat{\varphi}$$

$$\chi = 2g\varphi$$



Removing the string (1975)

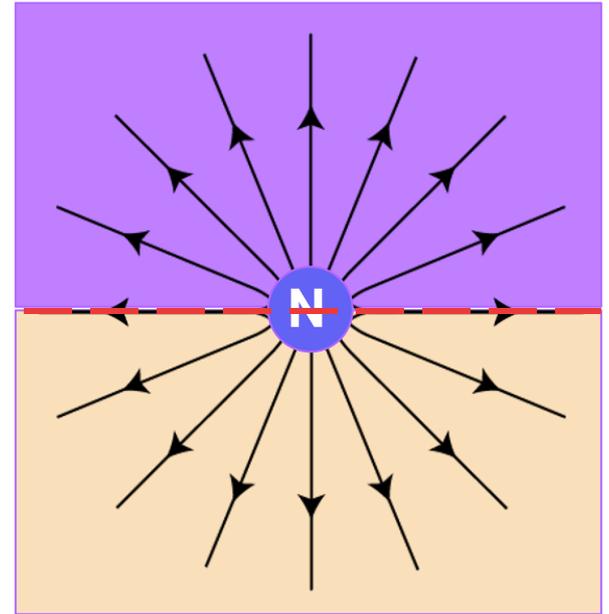
The same gauge field must be valid in QM:

$$U_g = \exp\left(\frac{-iq\chi(\varphi=0)}{\hbar c}\right) = \exp\left(\frac{-iq\chi(\varphi=2\pi)}{\hbar c}\right)$$

$$U_g = \exp(0) = \exp\left(\frac{-i4qg\pi}{\hbar c}\right)$$

$$\frac{2qg}{\hbar c} \in \mathbf{Z}$$

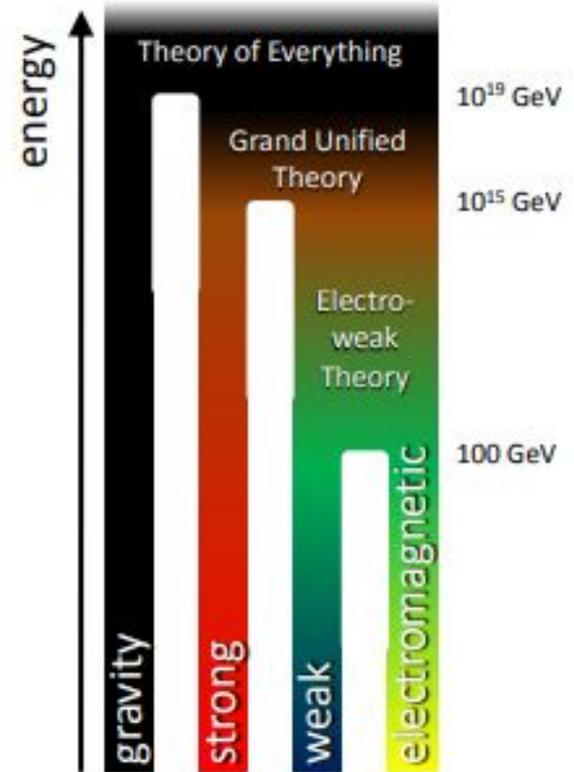
The charge quantization is a consequence of gauge symmetries!



'T Hooft-Polyakov Monopole(1975)

In simple terms

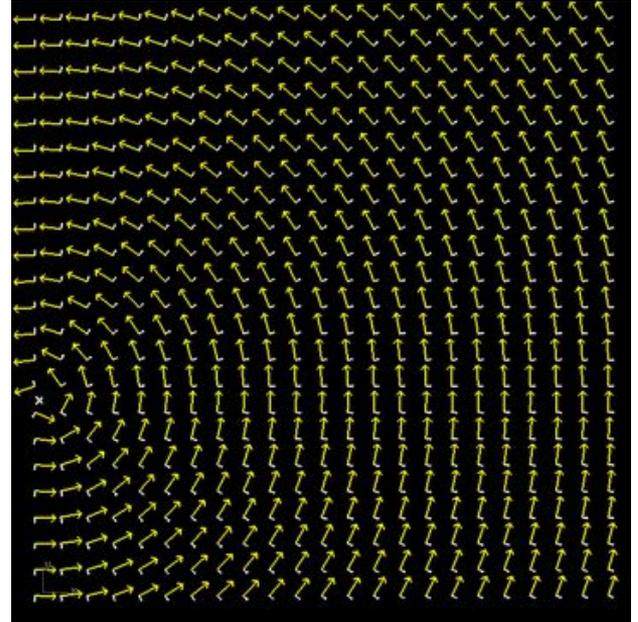
- The gauge field theory allows one to study the fundamental forces in a unified way in high energies.
- What separates the two symmetries, at lower energies, is the higgs field.



'T Hooft-Polyakov Monopole(1975)

In simple terms

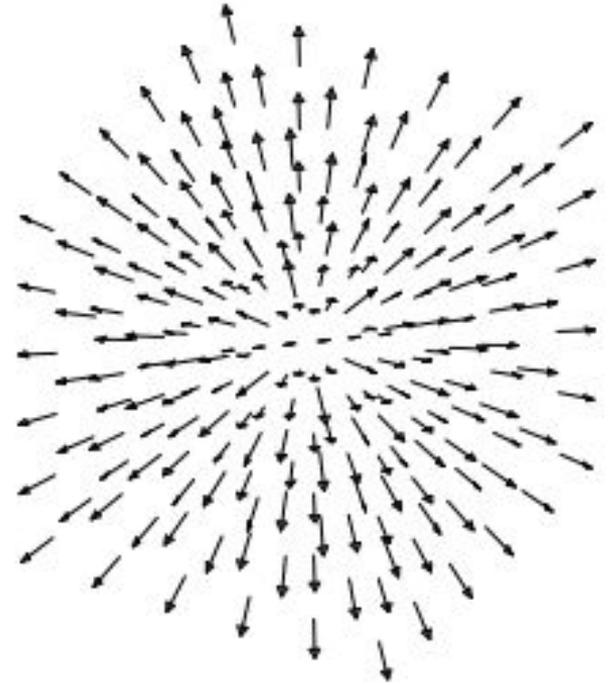
- The higgs field in Grand Unified Theories can be *represented as a vector field with two characteristics.*
 - Continuous and invariant under smooth changes (similar to vector field gauge invariance)
 - Non-zero everywhere in vacuum



'T Hooft-Polyakov Monopole(1975)

In simple terms

- The “Hedgehog” configuration is a interesting situation.
- Field is zero in the center, which can't be removed by smooth changes.
 - It cannot be vacuum: there is a massive particle in the center ($E \sim 10^{15} \text{ GeV}$)
- In this situation, the magnetic field can be determined as
$$\vec{B} = \frac{g}{r^2} \hat{r}$$
- Gauge theories inevitably predict magnetic monopoles!



2ST

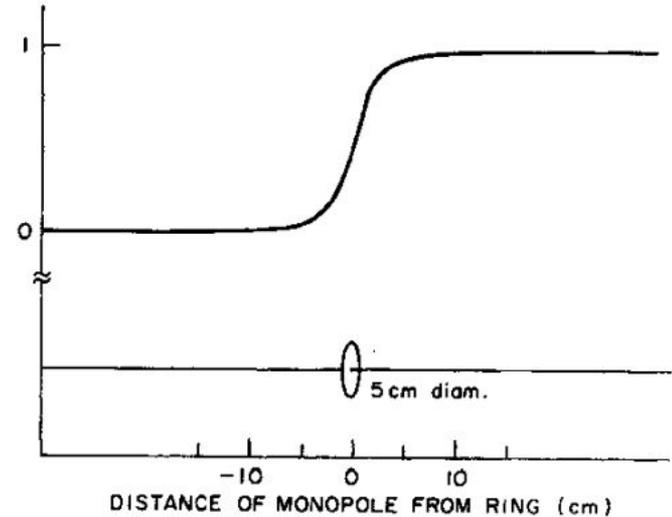
Looking for magnetic monopoles

Blas Cabrera's experiment

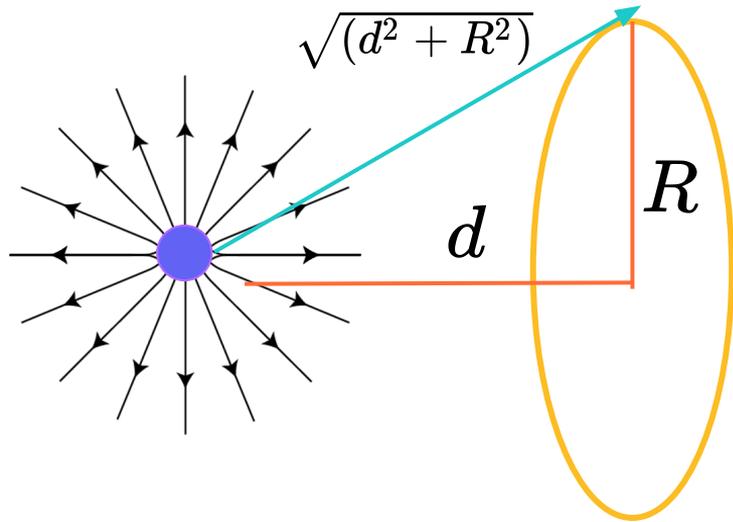
- Monopoles were never seen on earth, but is possible to look for them in cosmic radiation
- If a monopole passes through a superconductive coil, it will produce a current

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} - \frac{4\pi}{c} \vec{j}_m$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{1}{c} \frac{\partial \Phi_m}{\partial t} - \frac{4\pi}{c} \vec{I}_m$$



Blas Cabrera's experiment



- Taking $t=0$ as the time $d=0$:

- $t < 0$

$$\Phi_B = 2\pi g \left[1 + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$

- $t > 0$

$$\Phi_B = 2\pi g \left[-1 + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$

$$\Phi_B = \int_A \vec{B} \cdot d\vec{S}$$

$$\Phi_B = 2\pi g \left[1 - \frac{d}{\sqrt{(d^2 + R^2)}} \right]$$

$$\Phi_B = 2\pi g \left[1 - 2\Theta(t) + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$

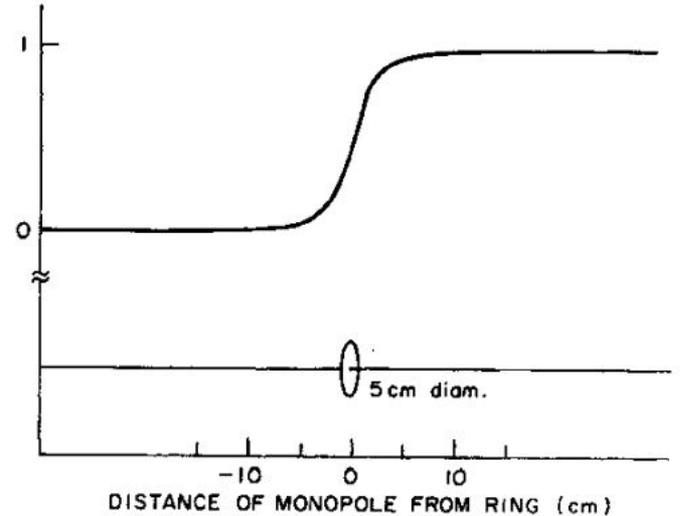
Blas Cabrera's experiment

$$-LI_e(t) = -\frac{1}{c}\Phi_m - \frac{4\pi}{c}g\Theta(t) \quad \Phi_B = 2\pi g\left[1 - 2\Theta(t) + \frac{vt}{\sqrt{(vt)^2 + R^2}}\right]$$

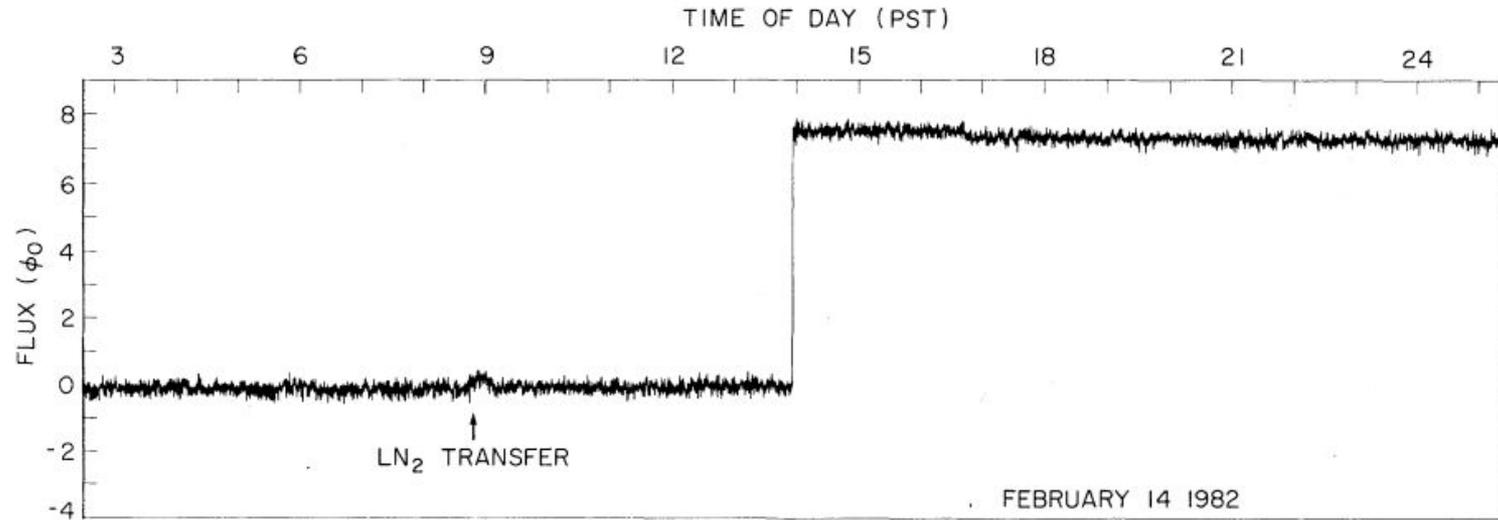
$$I(t) = \frac{2\pi g}{Lc} \left(1 + \frac{vt}{\sqrt{(vt)^2 + R^2}} \right)$$

$$\lim_{t \rightarrow +\infty} I(t) = \frac{4\pi g}{L}$$

$$\lim_{t \rightarrow -\infty} I(t) = 0$$



Blas Cabrera's experiment



Other experiments



Monopole, Astrophysics and Cosmic Ray Observatory (MACRO)

- **MACRO operated from 1989 to 2002 with a detection area of 10000m², never detecting one.**
- **No one ever saw any evidence of monopoles in space after 1982!**
- **If monopoles exist, would we be able to find them?**

The Poynting theorem with magnetic current

Poynting vector:

$$\vec{S} = \frac{c}{4\pi} \vec{E} \times \vec{B}$$

$$\nabla \cdot \vec{S} = -\frac{1}{4\pi} \left[\vec{B} \cdot \frac{\partial \vec{B}}{\partial t} + \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right] - \vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

Field's energy density:

$$u = \frac{E^2 + B^2}{8\pi}$$

$$\frac{\partial u}{\partial t} = \frac{1}{4\pi} \left[E \frac{\partial E}{\partial t} + B \frac{\partial B}{\partial t} \right]$$

$$\nabla \cdot \vec{S} + \frac{\partial u}{\partial t} = -\vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

The higher the currents, the faster the energy drops

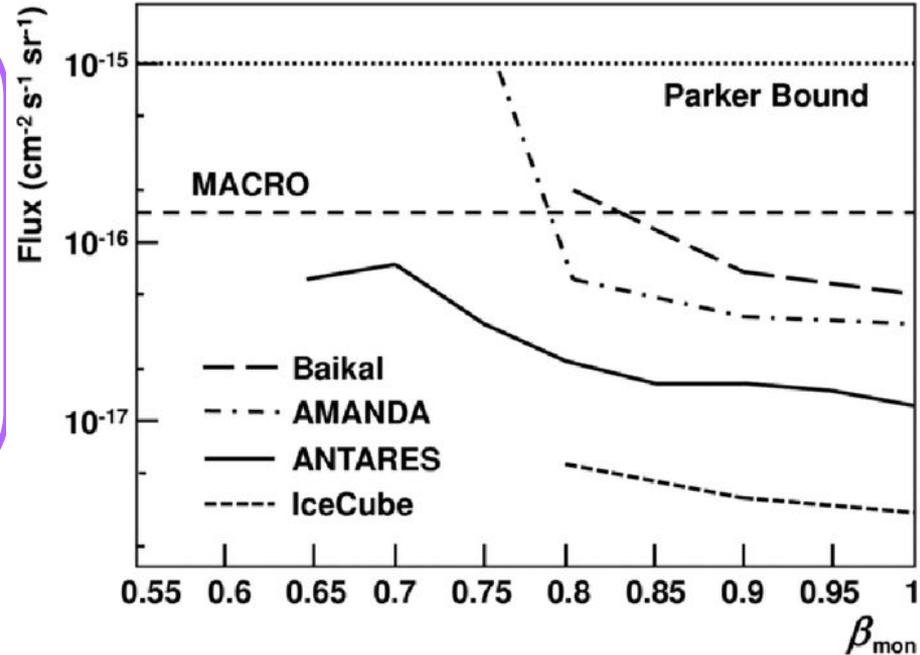
Astrophysical bounds

$$\nabla \cdot \vec{S} + \frac{\partial u}{\partial t} = -\vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

As the galaxy has a longing magnetic field of $3\mu G$ this drop cannot be so fast around here, so there is a limit to the magnetic current.

The flux F of magnetic field can be estimated. By Parker's bound, it would be

$$F \leq 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$



Astrophysical bounds

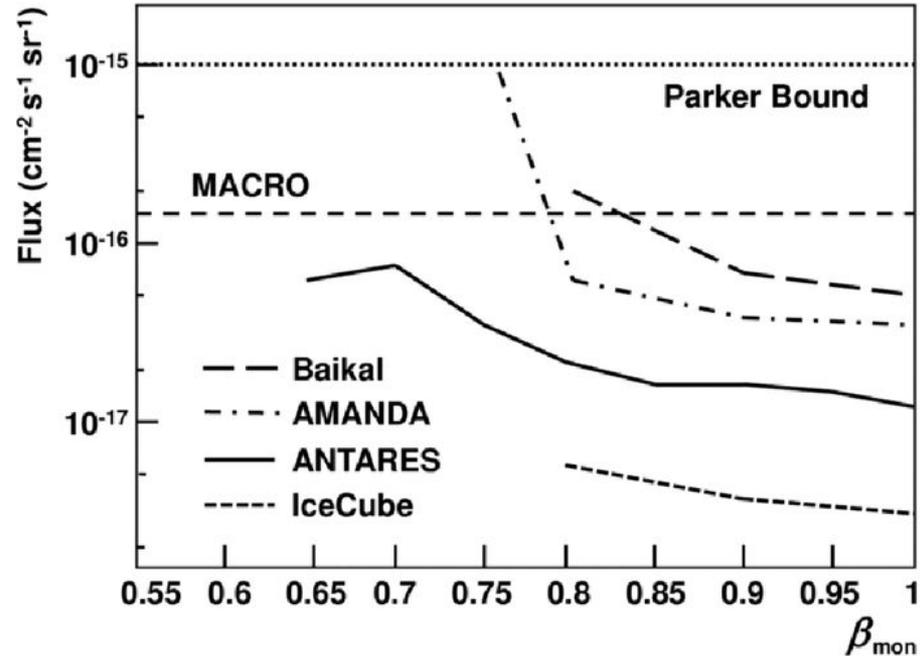
$$F \leq 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

Considering a detector the size of MACRO (10000m²), the rate of passing monopoles would be.

$$\frac{N}{t} \leq 1.5 \cdot 10^{-18} \text{ s}^{-1}$$

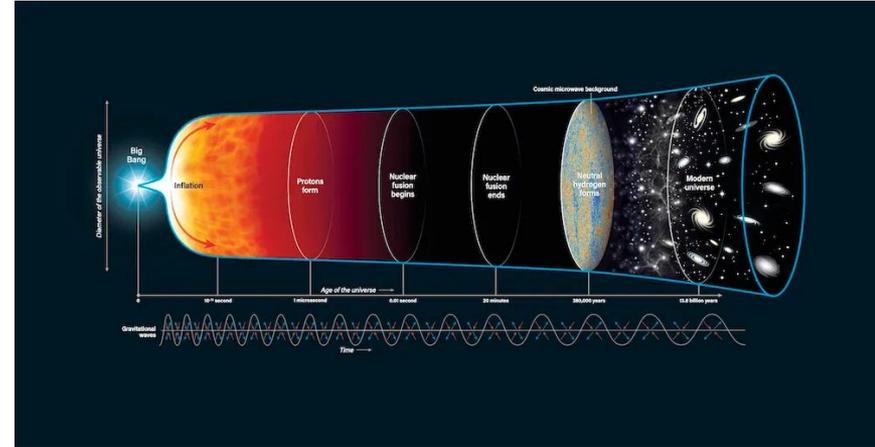
$$\frac{N}{t} \leq 5 \cdot 10^{-3} \text{ billion years}^{-1}$$

One every 20 billion years!?



Astrophysical bounds

- If they were as massive as expected ($E \sim 10^{15} \text{ GeV}$), they would have formed in the Big Bang.
- To compensate for their gravitational attraction, we need **cosmic inflation!** (“The monopole problem”)
- If that’s true, we will never find a monopole.
- Can we make one?

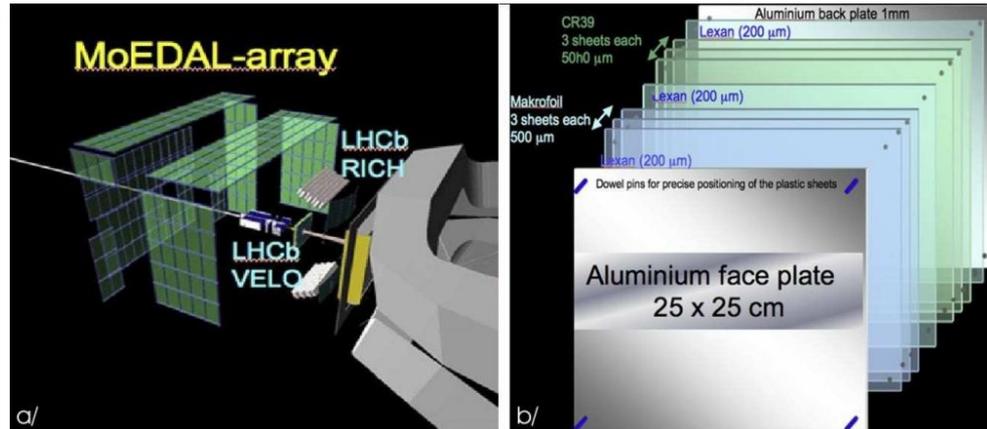


3ST

Creating a magnetic monopole

MoEDAL

- Although not predicted, Intermediate Mass Monopoles are still allowed.
- The Monopole and Exotics Detector at the LHC (MoEDAL), with collision energies of 8TeV, has been trying to create a monopole since 2009.



4ST

Conclusion

Conclusion

- **Magnetic monopoles are allowed under ED and QM**
- **Gauge field theories expect them at higher energies**
- **We never find one, and probably never will...**
- **MoEDAL is the most modern experiment currently trying to prove their existence**

Nobel prizes



Paul Dirac(1933)
Antimatter prediction



Carl Anderson(1936)
Discovery of the positron

References

- [1] Dirac, P. A. M. (1931). Quantised Singularities in the Electromagnetic Field. *Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences*, 133(821), 60–72.
- [2] Wu, T. T., Yang, C. N. (1975). Concept of nonintegrable phase factors and global formulation of gauge fields. *Physical Review D*, 12(12), 3845–3857.
- [3] Cabrera, B. (1982). First Results from a Superconductive Detector for Moving Magnetic Monopoles. *Physical Review Letters*, 48(20), 1378–1381.
- [4] Hooft, G. ' (1974). Magnetic monopoles in unified gauge theories. *Nuclear Physics B*, 79(2), 276–284. doi:10.1016/0550-3213(74)90486-6
- [5] Cabrera, B. (1982). First Results from a Superconductive Detector for Moving Magnetic Monopoles. *Physical Review Letters*, 48(20), 1378–1381.
- [6] Turner, M. S., Parker, E. N., Bogdan, T. J. (1982). Magnetic monopoles fields *Review D*, 26(6), 1296–1305.
- [7] Acharya, B., et al . . . Chatterjee, A. (2016). Search for magnetic monopoles with MoEDAL prototype trapping detector in 8 TeV proton-proton collisions at the LHC *Journal of High Energy Physics*, 2016(8).19

Thank you!

Special thanks: bruno Trebbi